

**THE GEOMETRY OF VECTOR BUNDLES AND  
AN INTRODUCTION TO GAUGE THEORY  
LECTURE 25**

PROFESSOR STEVEN BRADLOW  
CLASS NOTES FROM MATH 433

University of Illinois at Urbana-Champaign

March 20, 1998

During the last lecture we began to examine connections on tangent bundles,  $TM \rightarrow M$ . If  $M$  has dimension  $m$ , then this is a rank  $m$  bundle over  $M$ . Given a connection  $\nabla$  and local coordinates,  $(x_1, \dots, x_m)$  on  $M$ , we can produce a local frame,  $\{\frac{\partial}{\partial x_i}\}$  for  $TM$  and  $\{dx_i\}$  for  $T^*M$ . We saw that

$$\nabla\left(\frac{\partial}{\partial x_i}\right) = \Gamma_{ji}^k dx_k \otimes \frac{\partial}{\partial x_j}$$

This was how we defined the Christoffel symbols. The geodesic equation was derived from examining covariant constant vector fields. That is, given a curve  $\gamma(t)$  in  $M$ , we called  $\gamma(t)$  geodesic if the velocity vector field,  $\dot{\gamma}(t)$  was covariant constant,  $\nabla_{\dot{\gamma}}\dot{\gamma} = 0$ . In local coordinates,  $\gamma(t) = (x_1(t), \dots, x_m(t))$ , we get the following set of ordinary differential equations

$$\ddot{x}_k(t) + \Gamma_{ji}^k(\gamma(t))\dot{x}_i(t)\dot{x}_j(t) = 0, \quad k = 1, \dots, m$$

A consequence of the existence and uniqueness of solutions to differential equations implies that given any vector  $\vec{v} \in T_x M$ , there exists a unique geodesic curve,  $\gamma_{\vec{v}}(t)$  with  $\gamma_{\vec{v}}(0) = x$  and  $\dot{\gamma}_{\vec{v}}(0) = \vec{v}$ .

*Exercise 1.* Let  $\gamma_{\vec{v}}(t)$  be the geodesic through  $x$  in the direction of  $\vec{v}$ . If  $\lambda$  is some constant, show that  $\gamma_{\lambda\vec{v}}(t) = \gamma_{\vec{v}}(\lambda t)$ .

**Corollary.** Given a unit vector  $\vec{u} \in T_x M$  (a unit vector with respect to some Riemannian metric), for small enough  $\lambda$ ,  $\gamma_{\lambda\vec{u}}(t)$  will be defined at  $t = 1$ . We define a map,  $T_x M \rightarrow M$  defined on a small neighborhood of  $0 \in T_x M$  by  $\vec{v} \mapsto \gamma_{\vec{v}}(1)$ . This map is called the exponential map and is denoted by  $\exp(\vec{v})$ .

A fact from Riemannian geometry says that this is a diffeomorphism of the neighborhood of 0 in  $T_x M$  onto a neighborhood of  $x$  in  $M$ . The proof relies on the implicit function theorem. In order to set things up properly, we must examine the derivative of this map,  $D_0(\exp) : T_0(T_x M) \rightarrow T_x M$ .

*Exercise 2.* Show that  $D_0(\exp) = \text{Id}$ .

This defined geodesic coordinates.

**Claim.** With respect to these geodesic coordinates,  $\Gamma_{ij}^k$  vanishes at  $x$ .

**Sketch of Proof:** Given any  $\vec{v} \in T_x M$ , evaluate  $(\Gamma_{ij}^k dx_k)(\vec{v})$  using the geodesic along  $\vec{v}$ ,  $\gamma_{\vec{v}}(t)$ . Use the geodesic equation and note that in geodesic coordinates,  $\gamma_{\vec{v}}(t) = t\vec{v} = (x_1(t), \dots, x_m(t))$ .

## Torsion

**Definition 1.** For a connection,  $\nabla$ , on the tangent bundle,  $TM$ , over  $M$ , the *torsion* of the connection is a tensor field in

$$\Omega^0(\wedge^2(T^*M) \otimes TM)$$

defined by

$$\tau(X, Y) = \nabla_X Y - \nabla_Y X - [X, Y]$$

for any vector fields,  $X, Y \in TM$ .

It is not at all clear that  $\tau$  is actually an element of  $\Omega^0(\wedge^2(T^*M) \otimes TM)$ . The following exercise is in this direction.

*Exercise 3.* Show that  $\tau(X, Y)_b$  depends only on the values of  $X_b$  and  $Y_b$ . This will show that  $\tau_b \in \wedge^2 T_b^* M \otimes T_b M$ . Show that if  $f$  is a function, then  $\tau(fX, Y) = f\tau(X, Y)$ .

What does the torsion measure? With respect to the frames  $\{\frac{\partial}{\partial x_i}\}$  for  $TM$  and  $\{dx_i\}$  for  $T^*M$ , we see that

$$\begin{aligned} \tau\left(\frac{\partial}{\partial x_i}, \frac{\partial}{\partial x_j}\right) &= \nabla_{\frac{\partial}{\partial x_i}}\left(\frac{\partial}{\partial x_j}\right) - \nabla_{\frac{\partial}{\partial x_j}}\left(\frac{\partial}{\partial x_i}\right) \\ &= (\Gamma_{ji}^k - \Gamma_{ij}^k) \frac{\partial}{\partial x_k} \end{aligned}$$

So, if

$$\tau = \tau_{ij}^k dx_i \wedge dx_j \otimes \frac{\partial}{\partial x_k}$$

then  $\tau_{ij} = \Gamma_{ij}^k - \Gamma_{ji}^k$ . If  $\tau = 0$ , then the connection has symmetric Christoffel symbols.

**Fact:** Given  $\nabla$ , if  $\tau \neq 0$ , then we can modify  $\nabla$  to obtain a new connection,  $\tilde{\nabla}$ , with  $\tilde{\tau} = 0$ . The modification is given by the following procedure. We have

$$\Omega^1(\text{End}(TM)) = \Omega^0(T^*M \otimes \text{End}(TM))$$

and

$$T^*M \otimes \text{End}(TM) \cong T^*M \otimes (TM^* \otimes TM) \cong (T^*M \otimes TM^*) \otimes TM$$

Since  $\wedge^2(T^*M) \subseteq T^*M \otimes TM^*$ , we have that  $\tau \in \Omega^1(\text{End}(TM))$ . We can use this to define  $\tilde{\nabla} = \nabla - \frac{1}{2}\tau$ .

*Exercise 4.* Show that the torsion of  $\tilde{\nabla}$  is zero.

## The Levi-Civita Connection

Recall that a metric,  $g$ , on a manifold  $M$  if it is a bundle metric on the tangent bundle of  $M$ . We can thus ask for connections on  $TM \rightarrow M$  to be compatible with the metric. By our discussion of orthogonal connections, this can be expressed by the condition that  $\nabla g = 0$ .

**Claim.** Given  $g$ , there is a unique connection, called the Levi-Civita connection and denoted  $\nabla^{lc}$ , such that  $\nabla^{lc}g = 0$  and  $\nabla^{lc}$  is torsion free.

**Proof:** Use local coordinates  $\{x_1, \dots, x_n\}$ . Write

$$g = \sum g_{ij} dx_i \otimes dx_j$$

with respect to the local frame  $\{dx_i\}$  for  $T^*M$ . Set

$$\Gamma_{ij}^k = \frac{1}{2} g_{kl}^{-1} [g_{jl,i} - g_{ij,l} + g_{li,j}]$$

where  $g_{jli} = \frac{\partial}{\partial x_i}(g_{jl})$  to form the desired connection.  $\square$

We now get induced connections on  $T^*M, \otimes T^*M, \otimes TM$ , etc. On  $\wedge^p T^*M$ , we have  $\nabla^{lc} : \Omega^p(M) \rightarrow \Omega^0(T^*M \otimes \wedge^0 T^*M)$  and  $d : \Omega^p(M) \rightarrow \Omega^{p+1}(M)$ . We would like to relate these two maps.

**Claim.**  $Alt(\nabla^{lc}) = d$  where  $Alt$  is the alternation.

273 ALTGELD HALL, 1409 W. GREEN STREET, URBANA, IL 61801  
*E-mail address:* bradlow@math.uiuc.edu